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COMPARISON OF EXPERIMENTAL AND CALCULATED CRITICAL MASSES FOR A GAS CORE CAVITY REACTOR

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# COMPARISON OF EXPERIMENTAL AND CALCULATED CRITICAL MASSES FOR A GAS CORE CAVITY REACTOR

by Klaus H. Gumto
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## SUMMARY

This report presents the results of a comparison of calculated and experimental critical masses for a gas core cavity reactor. The comparison points out sources of error in the calculations and provides guidelines for further calculations.

This study uses the results of a series of experiments, conducted at the National Reactor Test Station, for comparison with analytical results. The calculations were performed on an IBM 7094 computer using both diffusion theory and transport theory codes. The transport theory calculations were made in cylindrical and spherical geometry, while the diffusion calculations were restricted to spherical geometry only. Input parameters included energy groups,  $P_0$  and  $P_1$  scatter-transfer cross sections, up and down scattering, number of materials and mesh interval spacing. The analysis covered the range of fuel radius to cavity radius ratios from 0.500 to 0.900.

The results indicate that the transport theory approximations used in this study overestimated the critical masses, while the diffusion calculations underestimated them. The deviation of the critical masses calculated with one-dimensional diffusion theory codes from the experimental critical masses was inversely proportional to the radius ratio. This points out the relative inability of diffusion theory to handle large voids. It also indicates that as the fuel radius is reduced, that is, as the radius ratio decreases, it becomes more difficult to approximate the fuel cylinder by a sphere. This problem also arose in the one-dimensional transport calculations, but it did not produce as large a deviation between the calculated and experimental critical masses as resulted from the diffusion calculations. Therefore, the one-dimensional transport calculations were more accurate at smaller radius ratio than the diffusion calculations.

The two-dimensional (cylindrical geometry) transport calculations had excessively long computer running times (several hours) and large critical mass deviations. These derivations were mainly due to lack of spatial and energy group definition imposed by computer storage limits. This problem did not arise in the one-dimensional calculations. Two-dimensional calculations will require optimization of input parameters and adequate core storage in order to get accurate results, although there is little prospect of significant reduction in computer time.

#### INTRODUCTION

The gas-core cavity reactor is potentially one of the most efficient means of space propulsion. The promise of high specific impulse has prompted numerous analytical studies of such power sources. On the other hand, experimental data on gas-core cavity reactors are scarce. To date, only two cavity reactors have been constructed and operated in the United States: one at Los Alamos and the other at the National Reactor Test Station (LPT). Both reactors simulated a true gaseous core by various means, and neither reactor was operated at the high temperatures which are expected in an operating gas core rocket reactor. Since previous analytical studies were based on mathematical approximations to theory, it was not certain that the results of these calculations were representative of an actual reactor.

In order to determine whether the critical mass of a gas-core cavity reactor could be calculated accurately, a series of calculations was performed at Lewis. Three configurations of the reactor at the National Reactor Test Station were used for the comparison. The calculations used the same geometry and materials as the experiment, with the exception that a homogeneous medium simulating gaseous uranium was assumed in the calculations instead of the  ${\rm UO}_9$  foils used to simulate gaseous fuel in the experiment.

The calculations were made in five distinct sets. Set 1, the initial calculations, used diffusion theory. In order to improve the results, transport theory was substituted. Sets 2 and 3 used a two-dimensional cylindrical geometry, while sets 4 and 5 had a one-dimensional, or spherical, geometry. These set numbers will be used throughout the report to identify each series of calculations.

Within the transport calculations, several parameters were changed from one set to another in order to approximate the experimental results more closely. These included the change in geometry, mesh spacing, number and spacing of energy groups and number of distinct regions. In addition, the effects of up and down scattering transfers and  $P_0$  against  $P_1$  cross sections were investigated. The comparison of the resulting critical masses with those obtained from the experiments provided a measure of the accuracy of the calculations and the effect of the changes. Henderson and Kunze (ref. 1) applied some of the conclusions of this study in their more detailed analysis of both the configurations in this report and those of later critical experiments with gaseous  $\mathrm{UF}_6$  fuel.

An IBM 7044/7094-II direct couple computer was used to perform all calculations. All computer runs were made with existing codes and cross section libraries.

#### **SYMBOLS**

D diameter

k multiplication factor

 $\boldsymbol{k}_{\mbox{eff}} \qquad \mbox{effective multiplication factor}$ 

L length

m fuel mass

P Legendre polynomial expansion term

#### CALCULATIONAL PROCEDURES

## Reactor Configuration

The reactor consisted of a cylindrical cavity 122 centimeters long and 183 centimeters in diameter, resulting in a cavity volume of 3.2×10 $^6$  cubic centimeters and a length-to-diameter ratio of 2/3. The cavity was formed by the inner walls of annular aluminum tank which surrounded the cavity with an 89-centimeter layer of  $D_2O$  as reflector and moderator. Aluminum structure weighing 764 kilograms was inside the tank. The tank was made in two parts, with one end on a movable table. This provided access to the cavity, but also introduced an air gap at the separation plane. The movable end reflector also contained a hole 30 centimeters in diameter to simulate an exhaust nozzle. However, this hole was closed with a  $D_2O$  filled plug for the configurations in this report. An aluminum support structure filled the cavity. This permitted the arrangement of the  $UO_2$  fuel sheets in a three-dimensional lattice which simulated gaseous fuel. These  $UO_2$  foils were 0.00254 centimeter thick, and were 93.2 percent enriched with  $U^{235}$  (ref. 2).

Three of the experimental fuel arrangements were selected for comparison. All consisted of foils so located that they formed a cylinder as long as, and concentric with, the cavity. Only the cylinder radius varied. The fuel radius to cavity radius ratios were 0.890, 0.667, and 0.530.

## Analytical Model

The analytical model is simplified from the experimental reactor in order to reduce the complexity of the calculations (fig. 1). This results in complete homogeneity

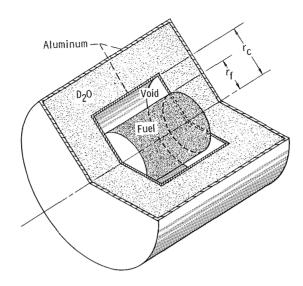


Figure 1. - Model of cavity reactor.

in each of the regions of the model reactor as compared to the discrete materials in the actual reactor. Also, the plugged exhaust nozzle and the air gap were eliminated from the model. The aluminum fuel support structure was homogenized over the entire cavity volume. For the model, this consisted of 34-kilogram aluminum with 0.51 kilogram of manganese impurity. The constituents of the  $\rm UO_2$  foils were homogenized over the fuel region which also contained part of the aluminum support structure. All materials were considered to be at room temperature for the calculations.

Since the experimental reactor was under construction at the time the analysis was started, the dimensions used were those currently available. Thus they will vary slightly from one set of calculations to another and from the published final dimensions of the experimental reactor. These variations are small, however, and do not significantly affect any calculations. Tables I to III give the dimensions used in each set of calculations.

In general, five regions were used

- (1) The fuel region, containing both fuel and support structure, with the fueled radius ranging from 0.5 to 0.89 of the cavity radius.
- (2) The void region, which is the annular space between the fuel and the cavity wall, containing only aluminum support structure.
- (3) The inner wall of the  ${\bf D_2}{\bf O}$  tank, which also forms the cavity and is made of aluminum.
- (4) The  $\rm D_2O$  with 0.25 percent  $\rm H_2O$  impurity and 693 kilograms of aluminum, all homogenized over the entire region.
  - (5) The outer tank wall, also made of aluminum.

# TABLE I. - MATERIALS AND GEOMETRY FOR ONE-DIMENSIONAL DIFFUSION

#### CALCULATIONS - SET 1

Distance to outer edge of region, cm	Number of mesh points	Width of mesh cell, cm	Material				
Fuel radius	to cavity rac	lius ratio, (	). 890				
81.30	35	2.323	Fuel				
91.34	. 5	2.007	Void				
92.86	3	. 507	Aluminum				
184.30	40	2.268	D <sub>2</sub> O				
185. 82	3	. 507	Aluminum				
Fuel radius	to cavity rad	lius ratio, (	), 667				
60.97	35	1.742	Fuel				
91.45	5	6.097	Void				
92.97	3	. 507	Aluminum				
184.44	40	2.268	D <sub>2</sub> O				
185.93	3	. 507	Aluminum				
Fuel radius	Fuel radius to cavity radius ratio, 0.530						
48.48	35	1.385	Fuel				
91.46	5	8. 596	Void				
92.98	3	. 507	Aluminum				
184. 45	40	2.268	$D_2O$				
185.94	3	. 507	Aluminum				

TABLE II. - GEOMETRY AND MATERIALS FOR TWO-DIMENSIONAL TRANSPORT CALCULATIONS

			1				
Distance to outer	Number of	Width of	Material				
edge of region,	mesh	mesh cell,					
cm	points	cm					
Set 2, radial direction							
<sup>a</sup> 81.38	<b>)</b> .	(a <sub>10.17</sub>	)				
<sup>b</sup> 60.96	8	b7.62	Fuel				
c <sub>48.43</sub>		c <sub>6.06</sub>					
	)	$a_{2.52}$	ĺ				
91.44	4	b7.63	Void				
		c <sub>10.74</sub>					
92.66	2	.61	Aluminum				
184.10	6	15.24	D <sub>2</sub> O				
185. 32	2	.61	Aluminum				
Se	t 2, axial di	ection					
1.22	2	0.61	Aluminum				
92.66	5	18, 29	D <sub>2</sub> O				
93.88	2	.61	Aluminum				
215.80	11	11.08	Fuel				
217.02	2	.61	Aluminum				
308.46	5	18.29	D <sub>2</sub> O				
309.68	2	.61	Aluminum				
Set	3, radial di	rection					
<sup>a</sup> 81. 38	)	(a <sub>10.17</sub>	1)				
<sup>b</sup> 60.96	8	b <sub>7.62</sub>	> Fuel				
c <sub>48.46</sub>		c <sub>6.06</sub>					
10, 10	)	$\binom{a_{2.52}}{}$	K				
91.44	4	b <sub>7.63</sub>	Void				
91, 77	T	c <sub>10.74</sub>					
185.32	10	9.39	D <sub>2</sub> O + aluminum				
Se	t 3, axial din	rection	1				
93. 88	9	10.43	D <sub>2</sub> O + aluminum				
215.80	11	11.08	Fuel				
309.68	9	10.43	D <sub>2</sub> O + aluminum				

<sup>&</sup>lt;sup>a</sup>Fuel radius to cavity radius ratio,  $r_*$ , 0.890. <sup>b</sup> $r_* = 0.667$ . <sup>c</sup> $r_* = 0.530$ .

TABLE III. - GEOMETRY AND MATERIALS FOR ONE-

#### DIMENSIONAL TRANSPORT CALCULATIONS

Fuel radius to cavity radius	Distance edge of		Numb me		l	th of	Material		
ratio	cı		poi		1	m			
	I	II	I	II	I	II			
	Set 4								
0.890	60.96	81.38	8	8	7.62	2.55	Fuel		
. 667	40.96	60.96	8	8	5. 12	2.50	Fuel		
. 530	40,96	48.46	8	8	5. 12	.94	Fuel		
	91.44		4				Void		
	103.44	185. 32	12	8	1.00	10.24	D <sub>2</sub> O + aluminum		
			Set 5	5					
0.900	70.00	82.30	8	4	8.75	3.08	Fuel		
. 890	73.00	81.38	1		9. 13	2.09	1		
. 875	70.00	80.01			8.75	2.50			
. 850	70.00	77.72			8.75	1.93			
. 825	70.00	75.44			8.75	1.36			
. 800	60.00	73. 15			7.50	3.29			
.775	60.00	70.87			7.50	2.72			
.750	60.00	68.58			7.50	2.15			
.700	60.00	64.00			7.50	1.00			
.667	56.00	60.69			7.00	1, 17			
.650	54.00	59.44			6.75	1.36			
.600	50.00	54.86			6.25	1, 22			
. 550	46.00	50.29			5.75	1.07			
.530	44.00	48.46			5.50	1. 12			
. 500	40,00	45, 72	🕴	♥	5.00	1.43	♥		
	91.44		4				Void		
	92.39		4		. 24		Aluminum		
	102.39	181.29	10	8	1.00	9.87	D <sub>2</sub> O		
	182.24		4		. 24	<del>-</del> -	Aluminum		

Since the calculations did not consider the nozzle or air gap present in the experiment, the critical masses chosen for comparison required a correction for 1.25 percent excess reactivity  $\Delta k$  due to their removal (ref. 1). This was added to the already existing excess reactivity in each configuration. A 3-percent impurity content in the  $UO_2$  foils was subtracted from the measured foil weight to get the uranium weight for the calculations.

Finally, a bias factor estimated by Henderson and Kunze (ref. 1) was applied in order to convert the critical foil mass to an equivalent critical mass for gaseous fuel.

TABLE IV. - EQUIVALENT CRITICAL MASS FOR GASEOUS FUEL

Fuel radius to	Actual fuel	Excess	Fuel worth,	Foil critical	Bias factor	Equivalent critical
cavity radius	loading,	reactivity,	∆k/k	mass at		mass for
ratio	kg U	Δk,	Δm/m	$k_{eff} = 1.00 \text{ kg U}$		gaseous fuel,
		percent		CII		kg U
0.890	10.7±0.1	1.46	0. 331	10.3±0.1	1.35±0.05	7.6±0.3
. 667	14.1±0.1	1.29	. 255	13.4 $\pm$ 0.1	1.21±0.05	11.1±0.5
. 530	19.8±0.2	2.23	. 182	17.4±0.2	1.13±0.05	15.4±0.8

The results of these steps are summarized in table IV.

The calculations are reported here in the order in which they were run. The first calculations used a one-dimensional diffusion theory code, which were followed by two-dimensional transport calculations in an effort to improve the results. When these did not achieve the accuracy desired due to computer storage limitations, the one-dimensional transport calculations were performed.

Following each set of calculations, the calculated  $k_{eff}$ , which were obtained by varying the fuel concentration, were plotted against their corresponding fuel masses, with the fuel radius to cavity radius ratio as a parameter. From these curves the critical masses at  $k_{eff}$  = 1.00 could be read for each radius ratio. These critical masses were then plotted against the radius ratio. This final graph permitted a comparison with another graph derived from experimental data.

#### Diffusion Calculations (Set 1)

The first calculations at Lewis were made with RP-4 a one-dimensional diffusion theory code. RP-4 uses 72 energy groups, including one thermal group. It obtains cross sections from its own library, has no up scattering and only one group scattering down.

For simplicity of calculations and to reduce the storage requirements of the computer code, a spherical geometry was used. Here the spherical dimensions were assumed to be equivalent to the radial dimensions of the cylindrical experiment configuration. The reactor was divided into the five regions which in turn were divided into 86 radial mesh intervals (see table I).

Two-dimensional diffusion calculations were not attempted. Hyland, Ragsdale, and Gunn (ref. 3) estimated that the computer time for an 80-centimeter radius cavity using only four energy groups would be 40 to 60 hours.

### Two-Dimensional Transport Calculations

A two-dimensional discrete angular segmentation (Sn) transport code programmed in FORTRAN IV (ref. 4), called TDSN, performed all transport theory calculations. Both the one- and two-dimensional calculations reported here were limited to the  $\mathbf{S}_4$  angular flux approximation which has been shown to be sufficiently accurate (ref. 5).

The first set of two-dimensional transport calculations (set 2) employed the same cross sections as the diffusion calculations (set 1). These had no up scattering and one group down scattering, and they were collapsed to seven broad energy groups from the original 72 groups as follows:

Group number	Energy range
1	7.79 MeV to 3.68 MeV
2	3.68 MeV to 1.31 MeV
3	1.31 MeV to 0.15 MeV
4	.15 MeV to 1.13 eV
5	1.13 eV to 0.414 eV
6	. 414 eV to 0. 125 eV
7	. 125 eV to 0 eV

The two-dimensional capability of TDSN permitted a cylindrical geometry with 22 radial and 29 axial mesh points encompassing five regions.

The second group of two-dimensional calculations (set 3) differed from the first group (set 2) in several respects. A major difference was in the cross sections. GAM-II (ref. 6) provided the fast cross sections, while GATHER-II (ref. 7) supplied the thermal cross sections. The cross sections of the constituents of each material were flux weighted over the composition of their respective regions. These were combined with MACGG, a code written at Lewis for this purpose. The resulting cross section set had up and down scattering and a  $P_0$  scatter transfer matrix. This is a set of scattering cross sections using a first-order approximation, that is, the  $P_0$  terms of a Legendre polynomial. A  $P_1$  scatter transfer matrix employs the second-order approximation, or the  $P_0$  and  $P_1$  terms of the Legendre polynominal (ref. 4). The cross sections were averaged over the following eight broad energy groups:

Group number	Energy range
1	14.90 MeV to 2.23 MeV
2	2.23 MeV to 0.82 MeV
3	.82 MeV to 7.10 keV
4	7.10 keV to 0.45 keV
5	.45 keV to 0.414 eV
6	. 414 eV to 0. 200 eV
7	.200 eV to 0.025 eV
8	.025 eV to 0 eV

This set again used a cylindrical geometry, also with 22 radial and 29 axial mesh points. These, however, included only three regions, that is, the fuel, void, and moderator regions. The moderator region also contained the inner and outer tank wall besides the  $D_2O$ . The calculations in sets 2 and 3 considered values of the fuel radius to cavity radius ratio of 0.890, 0.667, and 0.530.

## One-Dimensional Transport Calculations

The one-dimensional transport calculations also consisted of two sets. The first one (set 4) used the same cross sections, energy groups, and regions as set 3 of the two-dimensional calculations. The only change was to a spherical geometry with 40 radial meshpoints, and the fuel and  $D_2O$  regions were subdivided into two regions each. The inner and outer subregions are labeled I and II, respectively, in table III. These subregions permit closer mesh spacing at the fuel-void and  $D_2O$ -aluminum-void boundaries where the flux gradient is large. This results in an improved neutron flux calculation. The radius ratios for set 4 were 0.890, 0.667, and 0.530.

The cross sections in set 5 were also obtained from GAM-II and GATHER-II. These, however, were  $P_1$  cross sections. Again, a spherical geometry was used, this time with 42 radial meshpoints. The fuel and  $D_2O$  regions were subdivided to give the following 7 regions: (1), (2) fuel, (3) void, (4) inner aluminum tank, (5), (6)  $D_2O$ , and (7) outer aluminum tank. The radius ratios ranged from 0.500 to 0.900. The number of broad energy groups was increased to 20, with ten fast and ten thermal groups, having up scattering in the thermal groups and down scattering to all groups. The energy groups had the following ranges.

Group number	Energy range	Group number	Energy range
1	14.720 MeV to 3.012 MeV	11	2.382 eV to 1.600 eV
2	3.012 MeV to 2.019 MeV	12	1.600 eV to 1.200 eV
3	2.019 MeV to 1.225 MeV	13	1.200 eV to 0.850 eV
4	1.225 MeV to 0.608 MeV	14	0.850 eV to 0.330 eV
5	.608 MeV to 2.035 keV	15	0.330 eV to 0.220 eV
6	2.035 keV to 0.275 keV	16	0. 220 eV to 0. 100 eV
7	. 275 keV to 61. 440 eV	17	0. 100 eV to 0. 025 eV
8	61.440 eV to 10.680 eV	18	0.025 eV to 0.010 eV
9	10.680 eV to 3.928 eV	19	0.010 eV to 0.005 eV
10	3.928 eV to 2.382 eV	20	0.005 eV to 0 eV

## **RESULTS**

## Diffusion Calculations

The diffusion calculations produced critical masses lower than the experiment (table V, and figs. 2 and 3). This has been the experience at Lewis in other cavity reactor calculations using diffusion theory codes. When comparing the deviation of the calculated critical masses from the experimental critical masses, one can see that the deviation varies inversely with fuel radius to cavity radius ratio, or directly with the cavity void fraction. This result can be traced to two causes. First, the diffusion

TABLE V. - RESULTS OF DIFFUSION CALCULATIONS

Fuel radius to cavity radius ratio	Fuel mass, kg	Calculated, <sup>k</sup> eff	Extrapolated critical mass (k <sub>eff</sub> = 1.00), kg
0.890	8.9 9.8 10.2 10.7	1. 0666 1. 0919 1. 1027 1. 1142	7.1
0.667	9.7 10.7 12.0	1.0570 1.0774 1.0989	7.5
0. 530	10.8 11.7 13.2 14.8	1. 0239 1. 0331 1. 0496 1. 0640	8.9

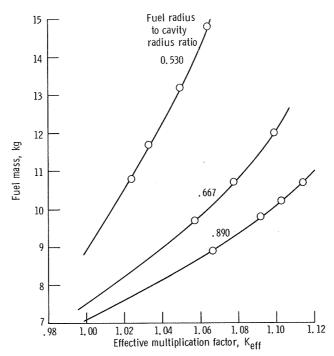


Figure 2. - Effective multiplication factor as function of fuel mass. One-dimensional diffusion calculations; 72 energy groups.

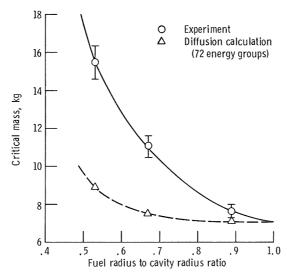


Figure 3. - Critical mass as function of fuel radius to cavity radius ratio. Diffusion calculation; experiment comparison.

theory codes cannot handle the large voids which exist at small fuel-radius to cavity-radius ratios. The second, a geometry problem, compounds the error.

Here the difficulty lies in trying to simulate a cylinder with a sphere. The problem arises in the fuel region, where the length to diameter ratio varies. Typical values for the length to diameter ratio are 3/4, 1/1, and 5/4 for radius ratios of 0.890, 0.667, and 0.530, respectively. A sphere can be expected to simulate a cylinder well if the length to diameter ratio is near 0.9; that is, the radius ratio is near 0.740. This geometry problem at small radius ratios also arose for the one-dimensional transport calculations.

Each diffusion calculation required 12 to 15 minutes of computer time.

## Two-Dimensional Transport Calculations

Both sets of the two-dimensional transport calculations resulted in critical masses 50 to 100 percent higher than the experimental critical masses (see table VI and figs. 4 to 6). The deviation increased rapidly for radius ratios smaller than 0.600.

TABLE VI. - RESULTS OF TWO-DIMENSIONAL
TRANSPORT CALCULATIONS

Fuel radius to cavity radius	Fuel mass,	Effective multiplication factor, $k_{\mbox{eff}}$		Critica kį	l mass,
ratio		Set 2	Set 3	Set 2	Set 3
0.890	10.0 10.5 11.0 11.5 11.8	0.9476  .9730 .9830 .9888 .9897	0.9566 .9660 .9747 	12.6	12.5
0.667	14.5 15.0 16.0 17.5 18.0 18.3	0.9762  .9927 1.0045 1.0072  1.0155	0.9569 .9682  .9876	16.9	20.2
0, 530	20.6 22.0 25.0 28.0	0.9657 .9725 .9845 .9935	0.9643	30.2	32+

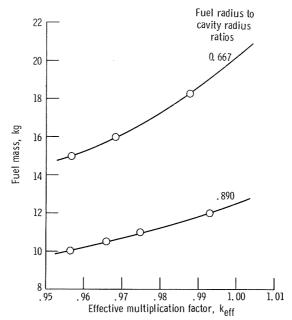


Figure 5. - Effective multiplication as function of fuel mass. Two-dimensional transport calculations; eight energy groups.

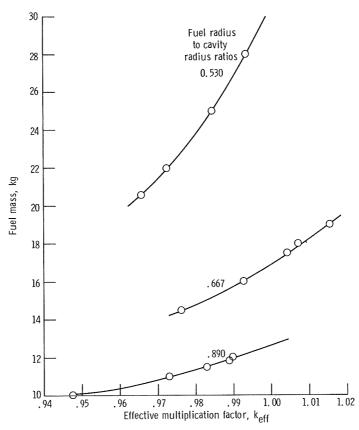


Figure 4. - Effective multiplication factor as function of fuel mass. Twodimensional transport calculations; seven energy groups.

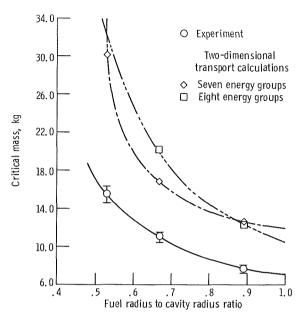


Figure 6. - Critical mass as function of fuel radius to cavity radius ratio. Two dimensional transport calculations; experiment comparison.

The addition of more groups, especially thermal groups, and up and down scattering improves the accuracy of the calculations. However, this is counteracted by the detrimental effect of reducing the number of regions and mesh points. Based on a comparison between sets 3 and 4, the overall result is to increase the error of the calculation, particularly for fuel-radius to cavity-radius ratios less than 0.850.

Henderson and Kunze (ref. 1) indicate that the mesh spacing in the reflector is the most critical. Here the large out-scattering cross sections result in very short mean free paths. These are about 2 centimeters for thermal neutrons. A rule of thumb for establishing mesh interval size is that it should be equal to or less than the mean free path in the material. All regions of every set meet this criterion except the  $D_2O$  region in sets 2 and 3. Here the mesh spacing is from 5 to 8 times the mean free path. This is the most likely source of error for the two-dimensional calculations.

Increasing the number of mesh points should produce greater accuracy, but, for a reactor as large as the one considered here, core storage requirements limit the number of mesh points permissible for two-dimensional calculations. The storage required is a function of the number of energy groups, materials, up and down scattering transfers, and geometry. At the storage limit, an increase in any of these quantities

requires a corresponding decrease in the number of one or more of the others. For any particular configuration, it should be possible to find an optimum combination of input parameters by trial and error, but this is a costly process since the computer time ranges from 150 to 900 minutes per calculation. It would probably be less costly and time consuming to determine the effect of changing each variable by doing one-dimensional calculations with TDSN before switching to two-dimensional calculations. Such a procedure, however, was not considered within the scope of this study.

### One-Dimensional Transport Calculations

The results of the one-dimensional calculations show that the critical masses again were overestimated (tables VII and VIII, figs. 7 to 11). The results (of the calculations in set 4) were more accurate than those of set 3, the two-dimensional calculations, which were identical except for cylindrical geometry and a smaller number of mesh points. This points out the importance of good spatial detail in the model. Again, the mean free paths were smaller than the mesh intervals for every region except the outer subregion of the  $D_2O$ . The same geometry problem described in the results of the diffusion calculations was also noted in these calculations. The greater accuracy in the calculations of set 5 is due to the greater number of energy groups, the thermal groups in particular, the use of the  $P_1$  scatter-transfer matrix, and better definition of mesh points at the fuel-void and  $D_2O$ -aluminum-void boundaries, where large flux gradients exist.

The computer program running times for set 4 calculations ranged from 15 to 30 minutes, while in set 5, the computer required from 28 to 32 minutes for a calculation.

TABLE VII. - RESULTS OF ONE-DIMENSIONAL

#### TRANSPORT CALCULATIONS

[Data for set 4.]

Fuel radius to	Fuel mass,	Effectively multiplication	Critical mass,
cavity radius	kg	factor,	kg
ratio		<sup>k</sup> eff	
0.890	9.80	0.9794	10.8
	10.70	.9979	] 10.0
0.667	8.95	0.8918	)
	10.60	. 9250	6.1
	12.10	. 9492	)
0. 530	11.80	0.8610	~27.0

TABLE VIII. - RESULTS OF ONE-DIMENSIONAL  $\label{transport} {\tt TRANSPORT\ CALCULATIONS\ -\ SET\ 5}$ 

Fuel radius to	Fuel mass,	Effective	Extrapolated
cavity radius	kg	multiplication	critical mass,
ratio		factor,	kg
musus de la companya		<sup>k</sup> eff	
0.900	9.25	1.0285	~8.4
. 890	11.00	1.0666	~8.4
. 875	11.00	1.0633	~8.5
. 850	9.25	1.0156	Ŋ
. 850	10.00	1.0366	8.7
. 850	11.00	1.0567	[]
. 825	10.00	1.0292	9.1
. 800	9.25	1.0000	
. 800	10.00	1.0225	9.3
. 800	11.00	1.0418	IJ
. 775	10.00	1.0123	9.5
. 750	11.00	1.0250	9.9
. 750	12.50	1.0381	
. 700	11.00	1.0060	)
.700	12.50	1.0162	10.9
. 700	13.50	1.0312	10.0
. 700	14.50	1.0435	
. 667	12.50	. 9950	100
. 667	13.50	1.0093	12.8
. 650	14.50	1.0153	ń
.650	17.50	1.0480	3.2
. 650	20.40	1.0679	J
.600	14.50	. 9809	]
.600	17.50	1.0120	
. 600	20.40	1.0302	16.2
.600	30.00	1.0700	IJ
. 550	20.40	. 9878	
. 550	30.00	1.0225	23.8
. 530	30.00	. 9983	30.2
. 500	20.40	. 9302	

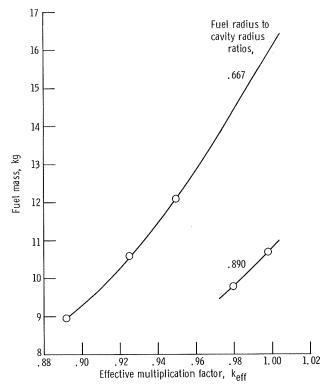


Figure 7. - Effective multiplication factor as function of fuel mass. One-dimensional transport calculations; eight energy groups.

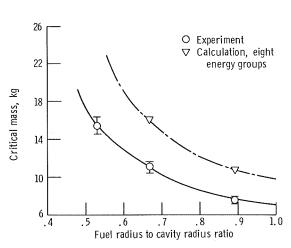


Figure 8. - Critical mass as function of fuel radius to cavity radius ratio. One-dimensional transport calculation; experiment comparison.

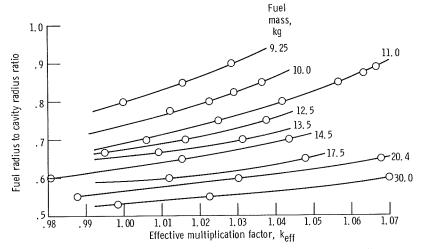


Figure 9. - Fuel radius to cavity radius ratio as function of effective multiplication factor. One-dimensional transport calculations; 20 energy groups.

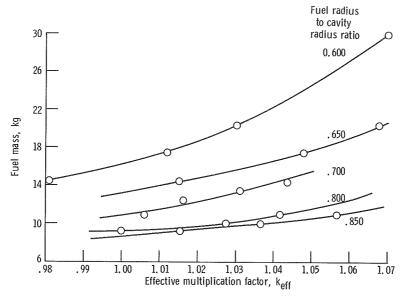


Figure 10. - Fuel mass as function of effective multiplication factor. One-dimensional transport calculations; 20 energy groups.

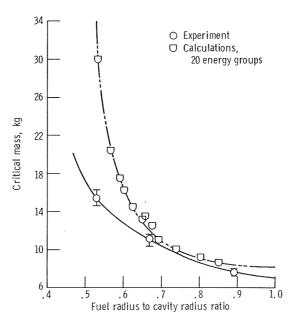


Figure 11. - Critical mass as function of fuel radius to cavity radius ratio. One-dimensional transport calculation; experiment comparison.

#### SUMMARY OF RESULTS

Calculated critical masses for a gas-core cavity reactor were compared with critical masses determined in an experiment. Three methods of calculations were used: (1) one-dimensional diffusion theory, (2) one-dimensional transport theory, and (3) two-dimensional transport theory. The mesh spacing, material distribution, and number of energy groups were varied among these to produce five distinct sets of calculations. These had the following results:

- 1. One-dimensional diffusion calculations underestimate critical masses. These calculations are good for fuel-radius-to-cavity-radius ratios near 1.0, but become increasingly inaccurate as the fuel becomes more compressed and the void fraction of the cavity increases. This makes them unsatisfactory for gas core cavity reactor calculations, as these reactors are expected to have large void fractions. These calculation required the least computer time, however.
- 2. Transport calculations overestimated critical masses. They provided better results than diffusion calculations at smaller radius ratios, although for radius ratios smaller than 0.600, the error increased rapidly for all types of calculations.
- 3. For computers with core storage and execution time limits, one-dimensional transport calculations produced better results than two-dimensional calculations. This is due to the greater refinement, that is, number of energy groups, mesh points, etc., possible with a one-dimensional geometry. The one-dimensional calculations also provided a decrease in computer time by a factor of 10 over the two-dimensional calculations.
- 4. The accuracy of transport calculations was improved by increasing the number of energy groups, thermal groups in particular, and by reducing the mesh spacing. Further gains in accuracy were achieved by a better definition of materials distribution, that is, adding regions, and by using a  $P_1$  scatter-transfer matrix instead of  $P_0$ .

Lewis Research Center,
National Aeronautics and Space Administration,
Cleveland, Ohio, May 8, 1970
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